

The Corrected Diffusion Approximation

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Summary

- What I do

I model light propagation in biological tissues for the purpose of locating early stage cancer cells.



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 - Early detection



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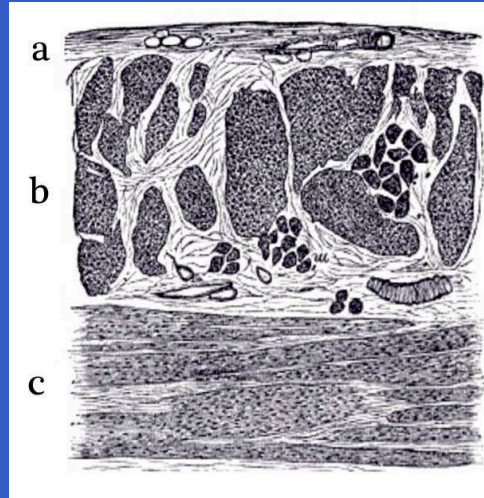
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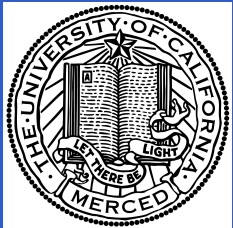
- Results



Tissue Structure



- a. Epithelial Layer
- b. Stromal Layer
- c. Smooth Muscle Layer



Setup

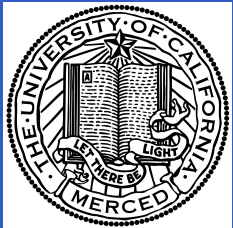
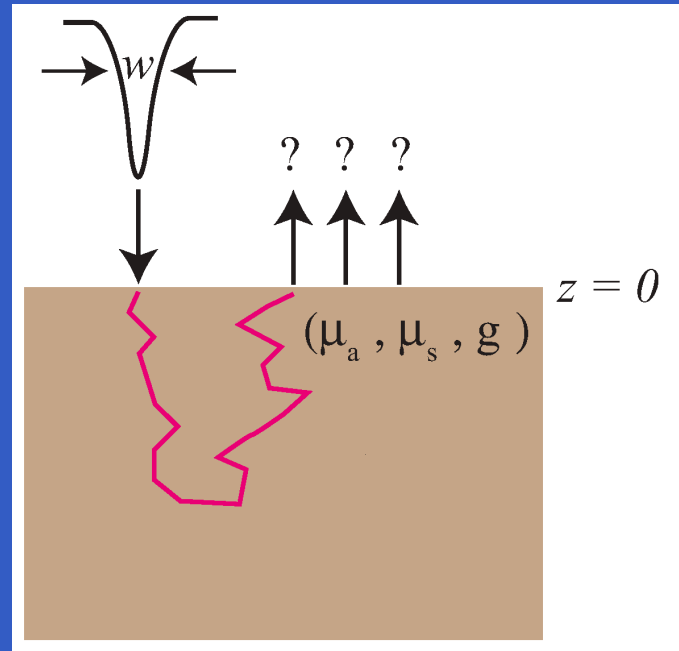
Goal: Model diffuse reflectance measurements of backscattered light by a turbid medium close to the source.



Setup

Goal: Model diffuse reflectance measurements of backscattered light by a turbid medium close to the source.

- Use a thin continuous beam incident normally on the medium
- Represent medium by a semi-infinite half space
- Given constant scattering and absorption coefficients



Light Propagation In Tissue

- Microscopic

Maxwell's equations provide a rigorous model for EM wave propagation



Light Propagation In Tissue

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Maxwell's equations provide a rigorous model for EM wave propagation

- Mesoscopic

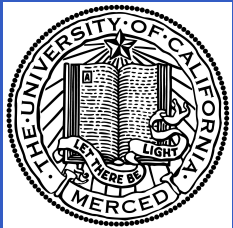
The Radiative Transport equation provides a model for light propagation as transport of particles



Radiative Transport Equation

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Radiative Transport Equation

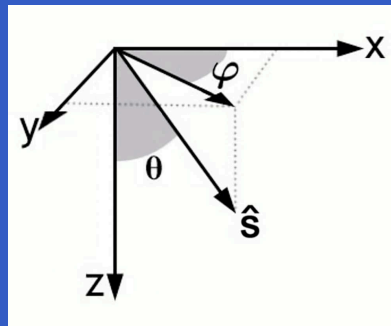
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p defines the fraction of power scattered in direction $\hat{\mathbf{s}}$ incident from direction $\hat{\mathbf{s}}'$.

$$\int_{\mathbb{S}^2} p(\hat{\mathbf{s}} \cdot \hat{\mathbf{s}}') d\hat{\mathbf{s}}' = 1$$



Boundary Condition

For a normally incident, Gaussian beam we consider

$$I(\mu, \varphi, x, y, 0) - r(\mu)I(-\mu, \varphi, x, y, 0) = \frac{\delta(\mu - 1)}{2\pi} f(x, y), \quad 0 < \mu \leq 1$$

$$f(x, y) = \frac{1}{2\pi w^2} e^{-\frac{x^2 + y^2}{2w^2}}$$



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and

$$I \rightarrow 0 \quad \text{as} \quad z \rightarrow \infty$$

In this, $r(\mu)$ is the Fresnel reflection coefficient at the boundary.



Light Propagation In Tissue

- Microscopic

 - Maxwell's equations provide a rigorous model for EM wave propagation

- Mesoscopic

 - The Radiative Transport equation provides a model for light propagation as transport of particles

- Macroscopic

 - The Diffusion Approximation is an approximation to the RTE.



Diffusion Approximation

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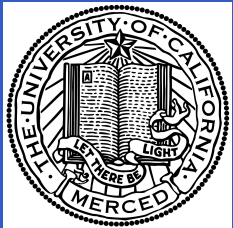


Diffusion Approximation

- We assume scattering is strong and absorption is weak ($\mu_s \gg \mu_a$)
- We assume isotropic scattering ($g = 0$)
- The Diffusion equation is of the form

$$\nabla \cdot (D \nabla \Phi) - \mu_a \Phi = S.$$

in this, $D = \frac{1}{3(\mu_a + \mu_s(1 - g))}$, and S is the interior source term.



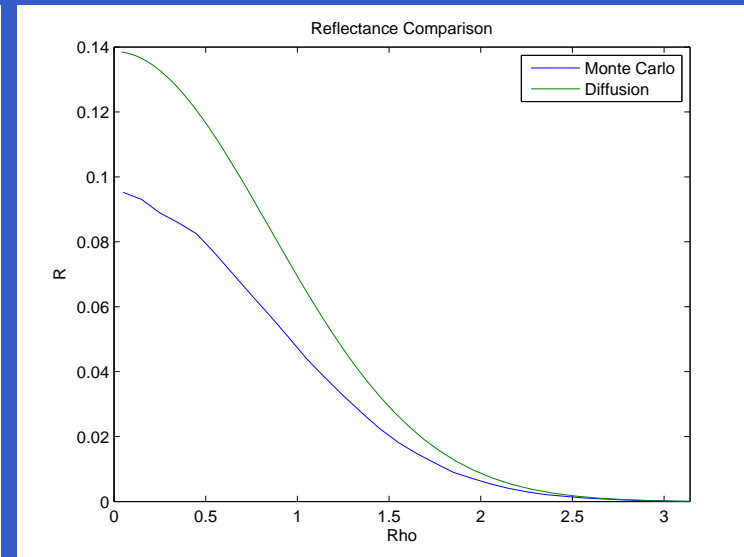
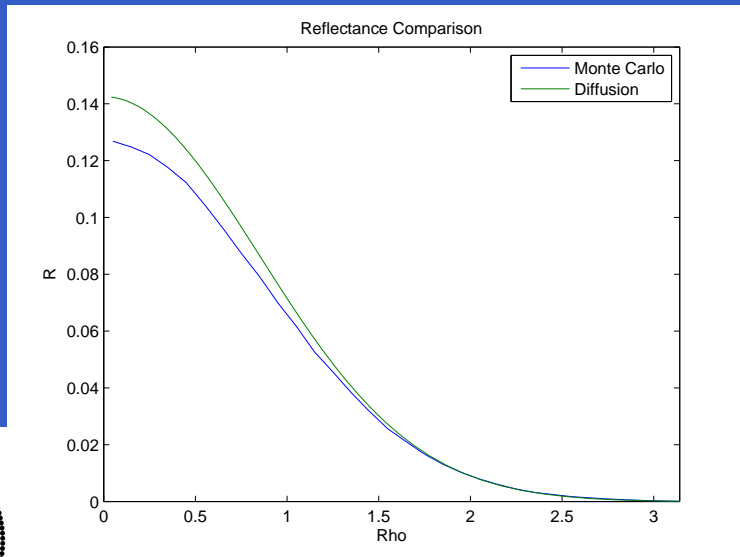
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Problem: This is known to be invalid close to the source.



Corrected Diffusion Model

Bridges the gap between Diffusion and RTE for tissues close to the boundary



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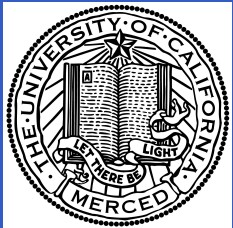
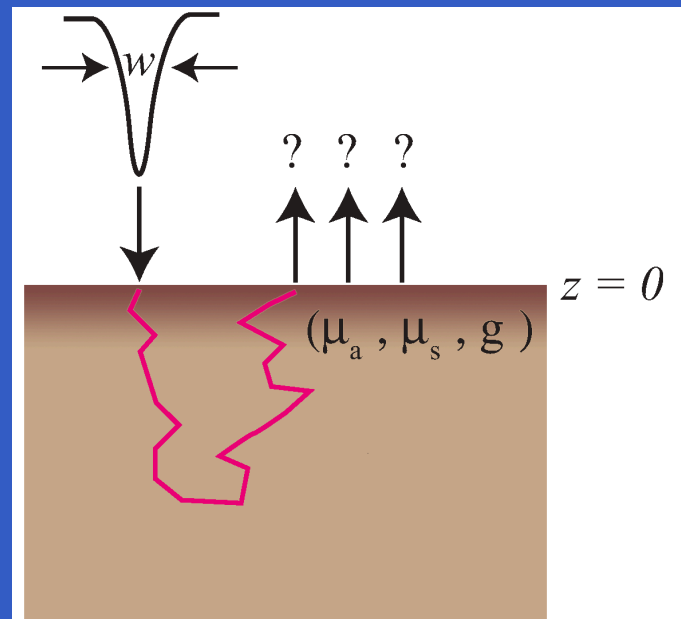


Corrected Diffusion Model

Bridges the gap between Diffusion and RTE for tissues close to the boundary

- Compute Interior Solution (Diffusion, Φ)
- Compute Boundary Layer Solution (RTE, Ψ)
- Combine results to satisfy original conditions

$$I(x, y, z, \hat{s}) = \Phi(x, y, z) + \Psi(x, y, z, \hat{s})$$



Derivation of CDA: Rescaling

- Three length scales in our analysis ($l_s \ll w \ll l_a$)

Scattering mean free path, $l_s = \frac{1}{\mu_s}$

Characteristic absorption length, $l_a = \frac{1}{\mu_a}$

Beam width w



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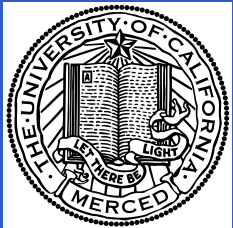
$$\text{Characteristic absorption length, } l_a = \frac{1}{\mu_a}$$

Beam width w

- Use our length scales to define small parameters α and β

$$\alpha = \frac{l_s}{l_a}$$

$$\beta = \frac{l_s}{w}$$



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- Use our length scales to define small parameters α and β

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- Rescale (x, y, z) with respect to w which nondimensionalizes the problem

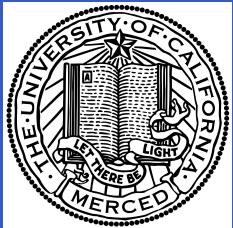
- Solve the rescaled, nondimensionalized equation using the fact that

$$\alpha \ll \beta \ll 1$$



CDA: Rescaled Problem

$$\beta\mu\partial_z I + \beta\sqrt{1-\mu^2}(\cos\varphi\partial_x I + \sin\varphi\partial_y I) + \alpha I + \mathcal{L}I = 0$$



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Subject to boundary conditions

$$I(\mu, \varphi, x, y, 0) = \frac{\delta(\mu - 1)}{2\pi} f(x, y) + r(\mu)I(-\mu, \varphi, x, y, 0), \quad 0 < \mu \leq 1,$$

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In these, $r(\mu)$ is the Fresnel reflection coefficient at the boundary.

We represent I as the sum of an interior solution and a boundary layer solution as in ^a.

$$(I = \Phi + \Psi)$$

^aS. B. Rohde and A. D. Kim, "Modeling the diffuse reflectance due to a narrow beam incident on a turbid medium," J. Opt. Soc. Am. A, **29**, 231-238 (2012).



Interior Solution

In solving for Φ , we find that ϕ_0 must satisfy the nondimensionalized diffusion equation

$$\nabla \cdot (\kappa \nabla \phi_0) - \frac{\alpha}{\beta^2} \phi_0 = 0.$$



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$$\Phi = \phi_0(\mathbf{r}) - \beta \hat{\mathbf{s}} \cdot [3\kappa \nabla \phi_0(\mathbf{r})] + O(\beta^2),$$



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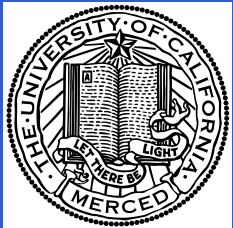
This solution alone cannot satisfy the boundary condition, and we apply a boundary layer solution



Boundary Layer Solution

Introduce stretched variable $z = \beta Z$, such that

$$\psi(\hat{\mathbf{s}}, x, y, Z) = \Psi(\hat{\mathbf{s}}, x, y, \beta Z),$$



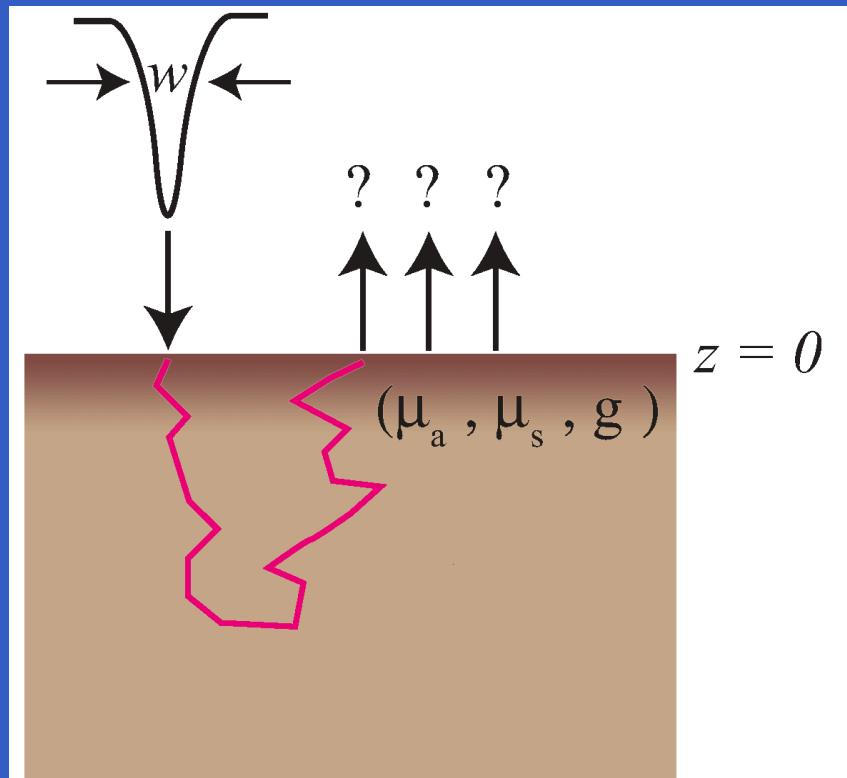
Boundary Layer Solution

Introduce stretched variable $z = \beta Z$, such that

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substitute into the RTE

$$\mu\psi_Z + \beta\sqrt{1-\mu^2}(\cos\varphi\psi_x + \sin\varphi\psi_y) + \alpha\psi + \mathcal{L}\psi = 0$$



Boundary Condition

We apply the modified boundary condition for $\psi = I - \Phi$

$$\psi(\mu, \varphi, x, y, 0) - r(\mu)\psi(-\mu, \varphi, x, y, 0) = \frac{\delta(\mu - 1)}{2\pi} f(x, y) - [1 - r(\mu)]\phi_0(x, y, 0) + 3\beta\kappa\mu[1 + r(\mu)]\phi_{0,z}(x, y, 0), \quad 0 < \mu \leq 1$$

Where $\psi = \psi_0 + \beta\psi_1$



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Where $\psi = \psi_0 + \beta\psi_1$, and ψ_0 satisfies the 1-D RTE

$$\mu\psi_{0,z} + \mathcal{L}\psi_0 = 0.$$

ψ_1 satisfies

$$\mu\psi_{1,z} + \mathcal{L}\psi_1 = -\sqrt{1 - \mu^2}(\cos \varphi\psi_{0,x} + \sin \varphi\psi_{0,y})$$



Asymptotic Matching

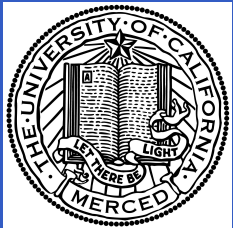
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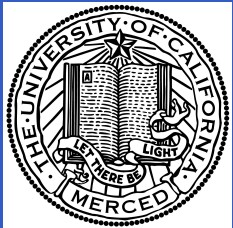
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We next solve for ϕ and then apply the full boundary condition with the numerically calculated Green's function to determine ψ

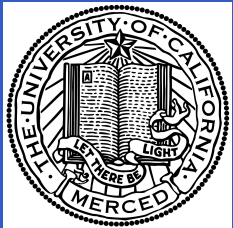


Interior Solution: Diffusion Approximation Solution

We can now solve

$$\nabla \cdot (\kappa \nabla \phi) - \alpha \phi = 0,$$

$$a_0 \phi - b_0 \phi_z = f_0 f(xy), \quad \text{at } z = 0.$$



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Using Fourier Transforms $(x, y) \rightarrow (\xi, \eta)$

$$-\xi^2 \kappa \hat{\phi} - \eta^2 \kappa \hat{\phi} + \kappa \partial_z^2 \hat{\phi} - \alpha \hat{\phi} = 0,$$



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$$-\xi^2 \kappa \hat{\phi} - \eta^2 \kappa \hat{\phi} + \kappa \partial_z^2 \hat{\phi} - \alpha \hat{\phi} = 0,$$

Since ϕ decays exponentially in z we set $\gamma(\xi, \eta) = -\sqrt{\alpha/\kappa + \xi^2 + \eta^2}$.

Substituting this into the BC we find

$$\hat{\phi} = \frac{f_0 \hat{f}(\xi, \eta)}{a_0 + b_0 \gamma}, \quad z = 0.$$



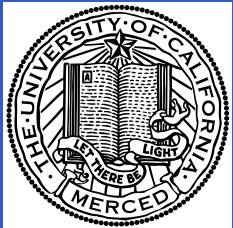
Reflectance Calculation

- Solve 1D RTE with Plane Wave Solutions



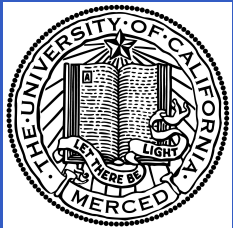
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Reflectance Calculation

- Solve 1D RTE with Plane Wave Solutions ^b
- Build Greens Function numerically
- Integrate with our source terms to solve for Ψ and Φ , $I = \Psi + \Phi$
- Integrate over the range of angles exiting the medium to determine reflectance at the boundary

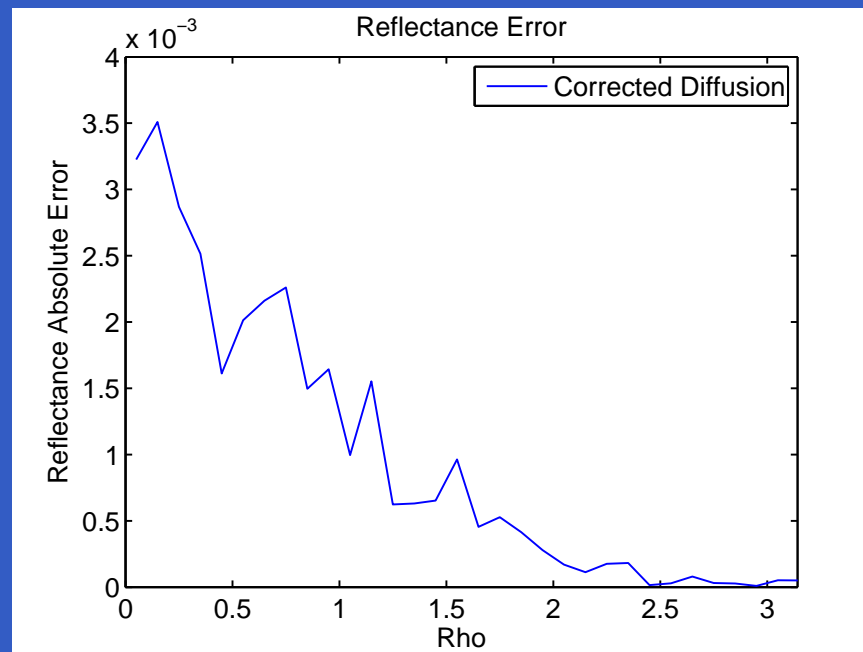
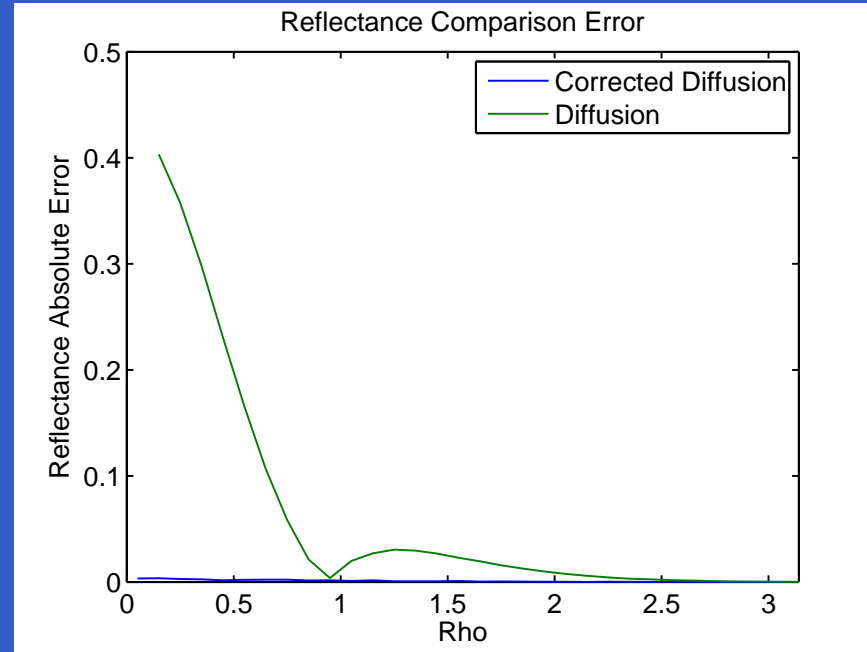
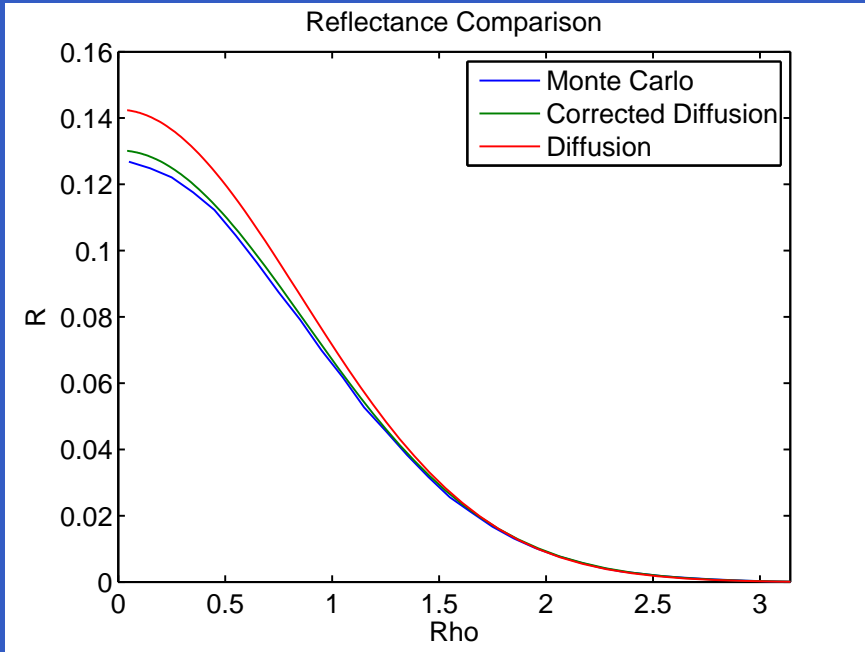
$$R(x, y) = - \iint_{NA} I(\mathbf{r}, \hat{\mathbf{s}}) \hat{\mathbf{s}} \cdot \hat{\mathbf{z}} d\hat{\mathbf{s}}.$$

^bA. D. Kim, “Correcting the diffusion approximation at the boundary,” J. Opt. Soc. Am. A **28**, 1007-1015 (2011).



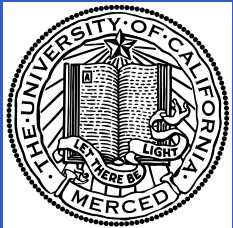
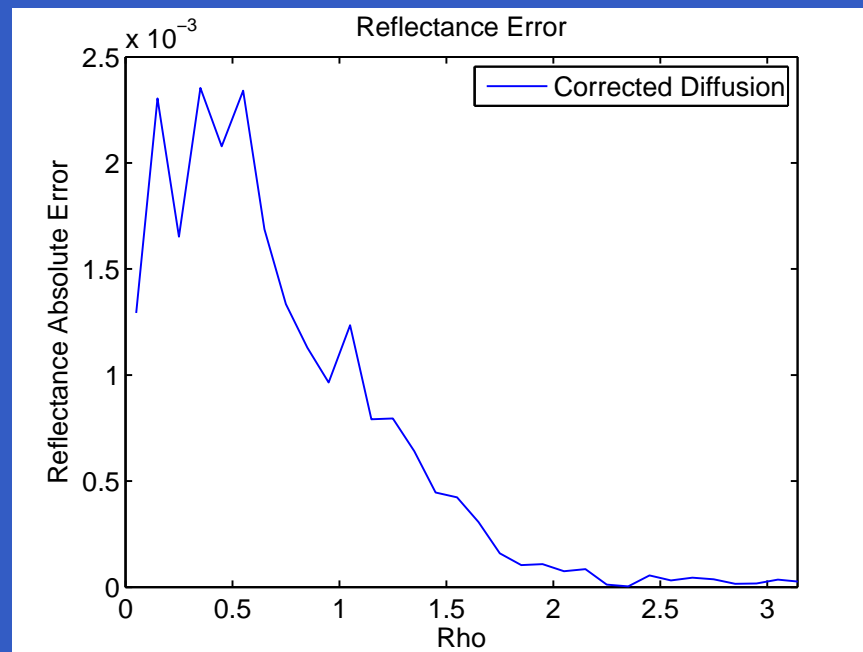
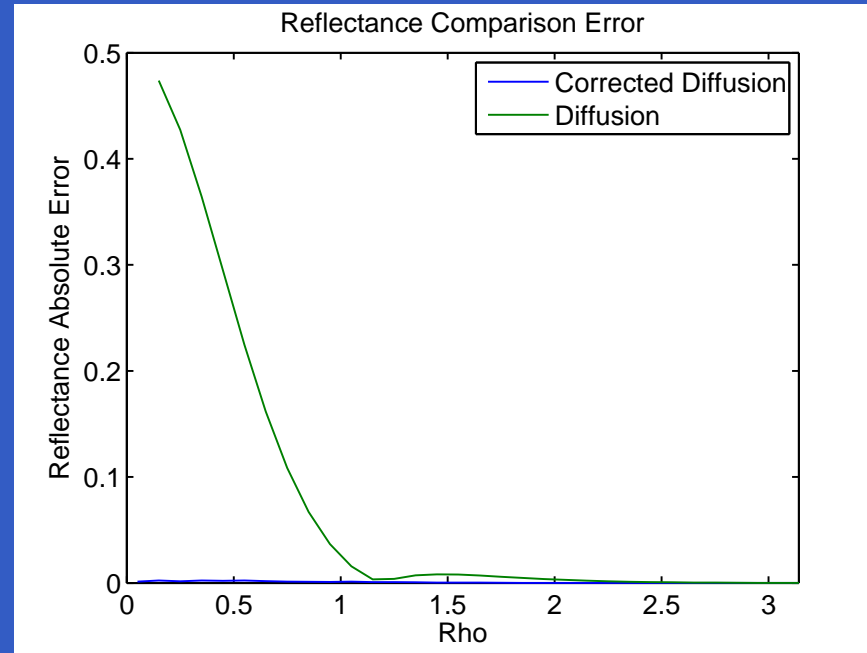
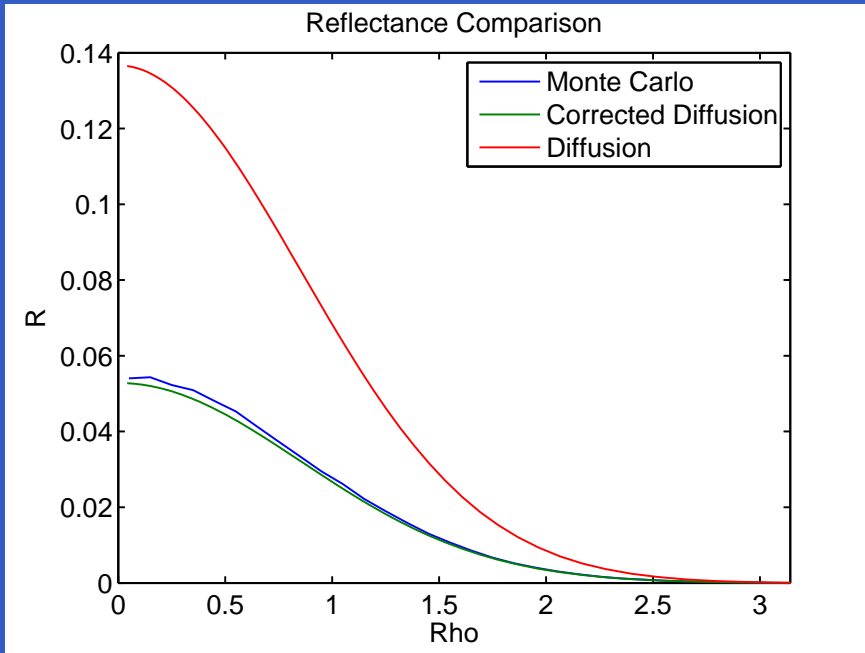
Results: How good is it?

$$\mu_a = .2(\text{mm})^{-1}, \mu_s = 100(\text{mm})^{-1}, g = 0.8, n_{rel} = 1.4, \text{BeamFWHM} = 1$$



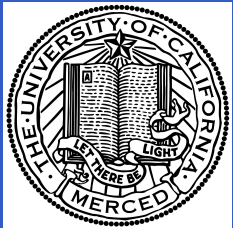
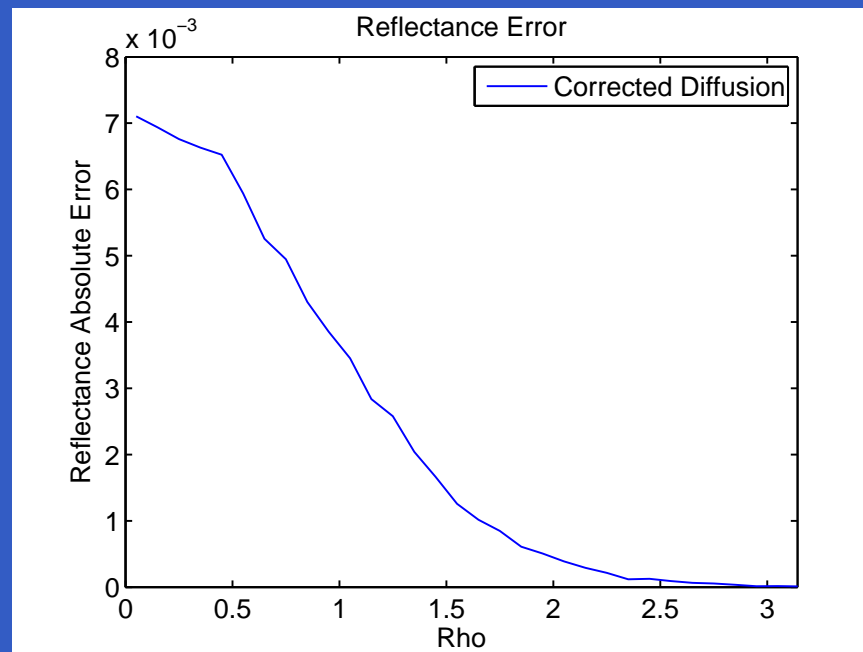
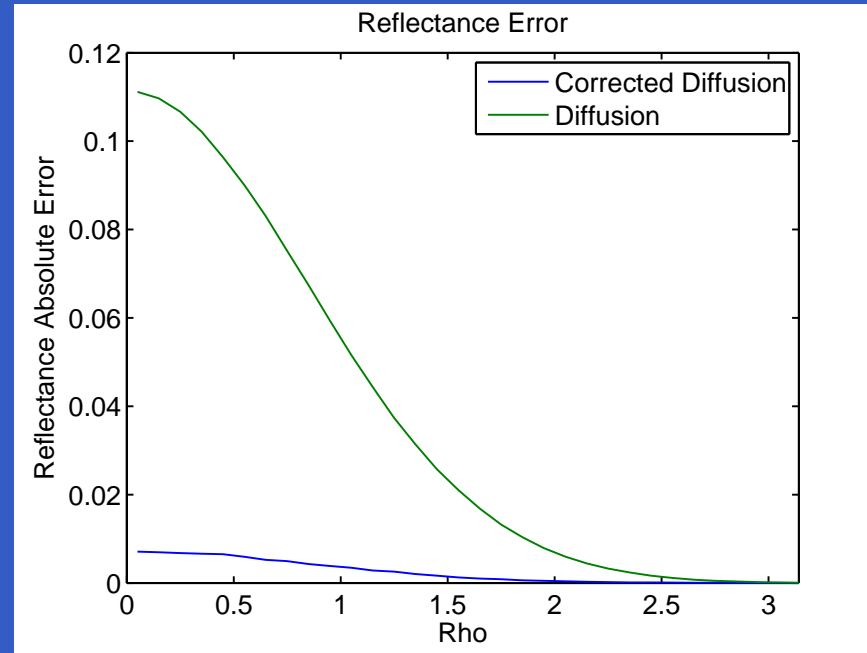
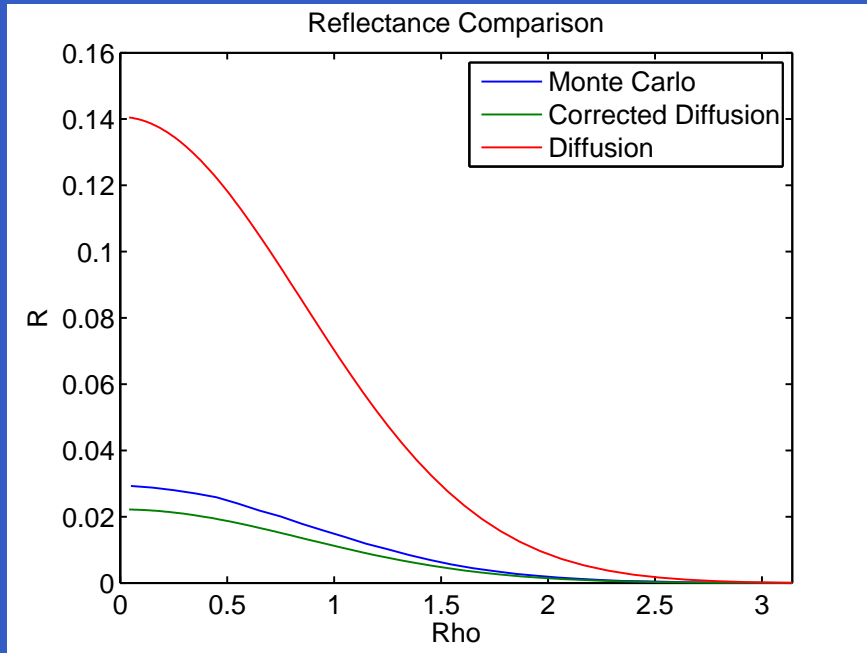
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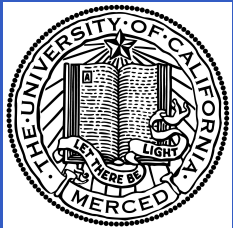
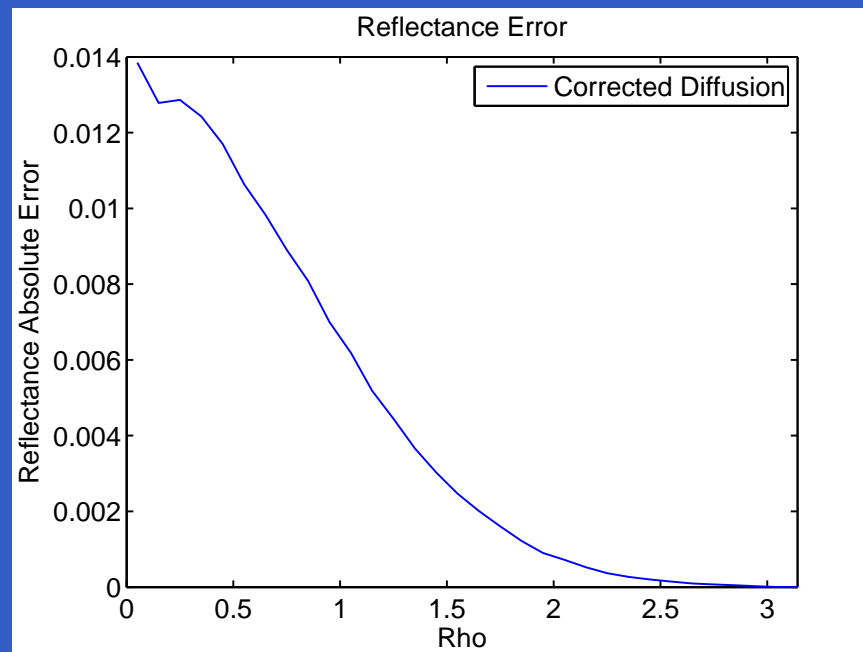
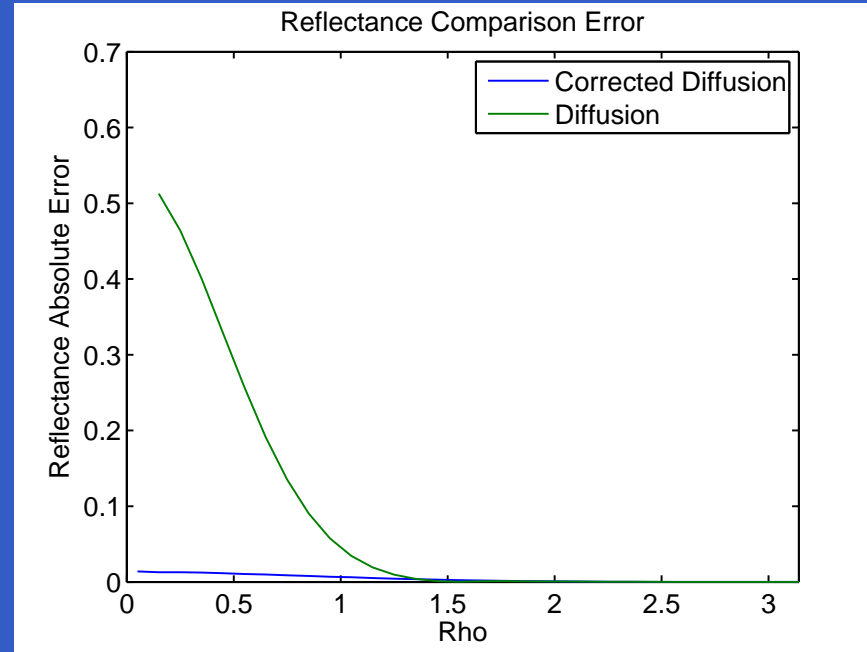
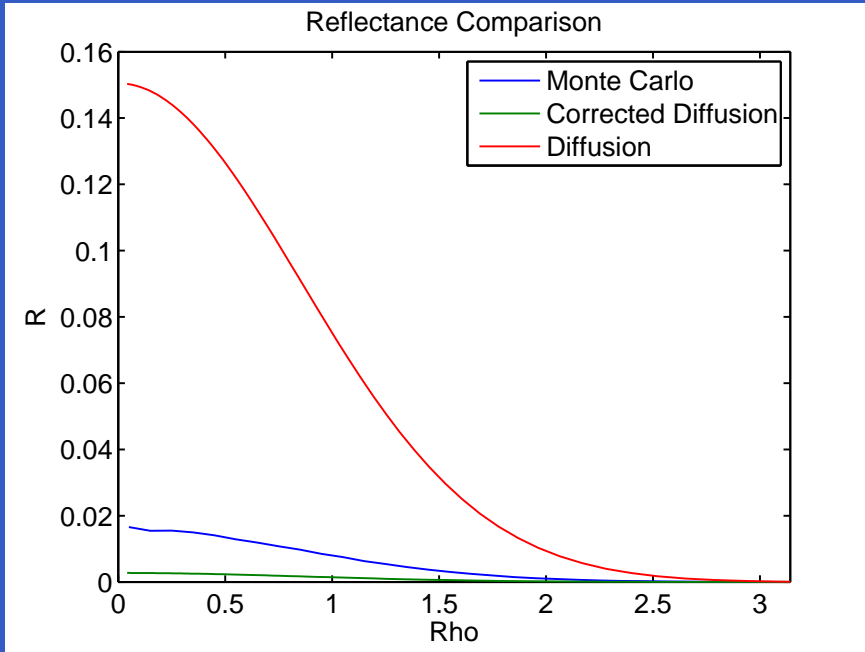
Results: How good is it?

$$\mu_a = 5(\text{mm})^{-1}, \mu_s = 100(\text{mm})^{-1}, g = 0.8, n_{rel} = 1.4, \text{BeamFWHM} = 1$$



Results: How good is it?

$\mu_a = 10(\text{mm})^{-1}$, $\mu_s = 100(\text{mm})^{-1}$, $g = 0.8$, $n_{rel} = 1.4$, $BeamFWHM = 1$



Conclusions and Acknowledgements

We constructed a forward model for accurate reflectance measurements close to the source

- We have extended it to include Fresnel reflection, layered tissues, and oblique incidence
- These models give us an option for modeling epithelial tissue specifically in an effort to locate early stage cancer cells, as well as an effective and invertible model for calculating optical properties of tissue



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To Do:

- Inverse problem
- Spatial frequency domain problem



Conclusions and Acknowledgements

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